

A DWBA analysis of the ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$ reaction
in the ${}^3\text{He}$ energy range 1.81 to 5.13 MeV

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Abstract

The angular distributions of the α - particles from the reaction ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$ measured in the ${}^3\text{He}$ energy range 1.81 to 5.13 MeV were analyzed in terms of the DWBA theory. Zero range DWBA was used without radial cut-off. Absolute spectroscopic factors were obtained and compared with the shell model calculations.

1. Introduction

The angular distributions of the α - particles from the ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$ reaction have been measured by different authors in the ${}^3\text{He}$ energy range 1.0 to 27.0 MeV [1 - 7]. DWBA calculations were carried out at ${}^3\text{He}$ energies between 5.0 and 18.0 MeV [6, 7]. In the present work, the experimental angular distributions of the α - particles from the ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$ reaction measured by P.D. Forsyth et al [2] at ${}^3\text{He}$ energies between 1.81 and 5.13 MeV were analyzed in terms of the DWBA theory.

2. DWBA analysis

DWBA calculations in the usual zero range approximation without radial cut-off were carried out on the Univac 1100 computer of the Institute of Nuclear Physics, Frankfurt, Deutschland-Germany using the code DWUCK due to Kunz [8]. The potential used has the form:

$$V(r) = V_0(r) + V f(x) + i W f(x_1) - V_a \frac{1}{r} \frac{d}{dr} f(x) \vec{L} \cdot \vec{S} \quad (1)$$

$$\text{where } f(x) = \left[1 + \exp \frac{(r - r_0 A^{1/3})}{a} \right]^{-1}$$

$$f(x_1) = \left[1 + \exp \frac{(r - r_1 A^{1/3})}{a_1} \right]^{-1}$$

The Coulomb potential was assumed to be that of a uniformly charged sphere of radius $R_c = r_c A^{1/3}$. The neutron was assumed to be captured in a potential well of the Saxon - Woods form with the n-pi-orbit potential of the Thomas type. The well depth parameter was adjusted by the code to give the proper neutron binding energy. The optical model parameters are collected in Table 1. All the ^3He and the α -particle parameters were tried in the present analysis. It is found however that the results of the analysis are sensitive to the choice of the optical model parameters. Satisfactory fits with reasonable spectroscopic factors were obtained with optical model parameter set BC deduced from the general optical model potentials present in ref. [9]. The results of the analysis are shown in figs. 1, 2, 3. The spectroscopic factors were calculated by using the formula:

$$C^2 S = \frac{1}{N^2} \frac{2(2j+1)}{2s+1} \frac{\sigma_{\text{exp.}}(\text{peak})}{\sigma_{\text{DW}}(\text{peak})} \quad (2)$$

where s is the spin of the captured neutron,
 j is the total angular momentum of the captured neutron,
 $\sigma_{\text{exp.}}$ is the experimental reaction cross-section at the peak,
 σ_{DW} is the calculated cross-section at the peak,

is a normalization factor and its value in the case of (${}^3\text{He}, \alpha$) reactions = 25.0

$C^2 = (T_0 t, M_T, M_t \quad T M_T)^2$ is the square of the Clebsch-Gordan coefficient for isobaric-spin coupling [10], where T_0, T and t are the isobaric-spin quantum numbers of the ${}^6\text{Li}$ final state, ${}^7\text{Li}$ initial state and of the transferred neutron respectively, and M_T, M_T, M_t are their Z -components. In the case of the reaction ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$, $C^2 = 1$ for $T = 0$ states and $C^2 = 1/3$ for $T = 1$ states.

The spectroscopic factors are given in Table 2. Theoretical spectroscopic factors have been calculated by S. Cohen et al [14] and by S. Varma et al [15]. S. Cohen et al listed separate components for $1p_{3/2}$ and $1p_{1/2}$ transfers, but S. Varma et al gave only the sum. The experimental spectroscopic factors for $1p$ transfers obtained in the present work were evaluated as though they arose entirely from $1p_{3/2}$ transfers or, if allowed entirely from $1p_{1/2}$ transfers. The spectroscopic factors obtained in the present work are generally higher than the theoretical predictions.

Conclusions

The results of the present analysis are sensitive to the choice of the optical model parameters. However, the general formulae of the optical model potentials given in ref. [9] seem to be applicable to the case of the interactions of ${}^3\text{He}$ with light nuclei at low energies [11]. The fits obtained in the present work for the different α -groups from the ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$ reaction at the ${}^3\text{He}$ energies under consideration are generally good at forward angles. The lack of obtaining better fits in the case of the α_0 and α_{2^-} groups at backward angles is attributed to an exchange mechanism in which the α -particle comes directly from the target nucleus, the ${}^3\text{He}$ being captured [3], while the lack of obtaining better fits in the case of the α_1^- group may be attributed to the compound nucleus effects.

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Fig.1 DWBA fits to the angular distributions of the α_0 - group at the ${}^3\text{He}$ energies 4.36 and 5.13 MeV. Dots are the experimental results taken from ref. [2]. Solid lines are the DWBA calculations in the usual zero range approximation without radial cut-off using the optical model parameter set BC present in Table 1.

Fig.2 DWBA fits to the angular distributions of the α_1 - group from the reaction ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$ at ${}^3\text{He}$ energies between 1.81 and 2.98 MeV. Dots are the experimental results taken from ref. [2]. Solid lines are the DWBA calculations in the usual zero range approximation without radial cut-off using the optical model parameter set BC present in Table 1.

Fig.3 DWBA fits to the angular distribution of the α_2 - group from the reaction ${}^7\text{Li}({}^3\text{He}, \alpha){}^6\text{Li}$ at $E_{{}^3\text{He}} = 4.36$ MeV. Dots are the experimental results taken from ref. [2]. Solid line represent the DWBA calculations in the usual zero range approximation without radial cut-off using the optical model parameter set BC present in Table 1.

Reaction channel	Ref.	V (MeV)	r_0^V (fm.)	a_0^V (fm.)	r_1^V (fm.)	a_1^V (fm.)	V_S (MeV)	r_0^S (fm.)	a_0^S (fm.)	r_1^S (fm.)	a_1^S (fm.)	r_0^C (fm.)
A Helium-3	[12]	140.0	4.30	0.67	7.7	0.75	12.0	1.20	0.70	1.3		
B Helium-3	[9]	34.0	1.50	0.65	4.0	1.50						1.6
C Helium-4	[9]	40.0	1.70	0.65	12.0	1.70						1.7
D Helium-4	[13]	189.4	1.92	0.60	7.7	1.92	11.1	1.92	0.60			1.4

Table 2 Spectroscopic factors

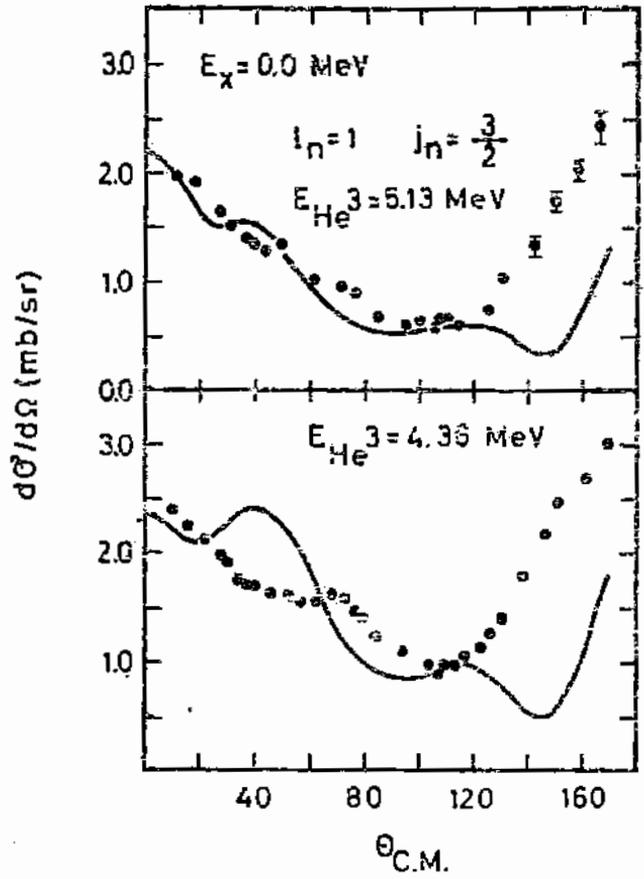
E_{He3} (MeV)	S exp. (pres. work)				S th. [14]		S th. [15]		
	$\alpha_1(a)$	$\alpha_1(b)$	$\alpha_2(c)$	$\alpha_0(d)$	α_1	α_2	α_1	α_2	
	$J_n=3/2$	$J_n=1/2$	$J_n=3/2$	$J_n=3/2$	$J_n=3/2$	$J_n=1/2$	$J_n=3/2$	$J_n=3/2$	
5.13	1.07 or 1.14			0.43	0.29				
4.36			1.12			0.89		0.75	
≈ 2.98		1.37			0.55			0.35	
2.50		1.29							
2.14		1.04							
1.81		1.12							

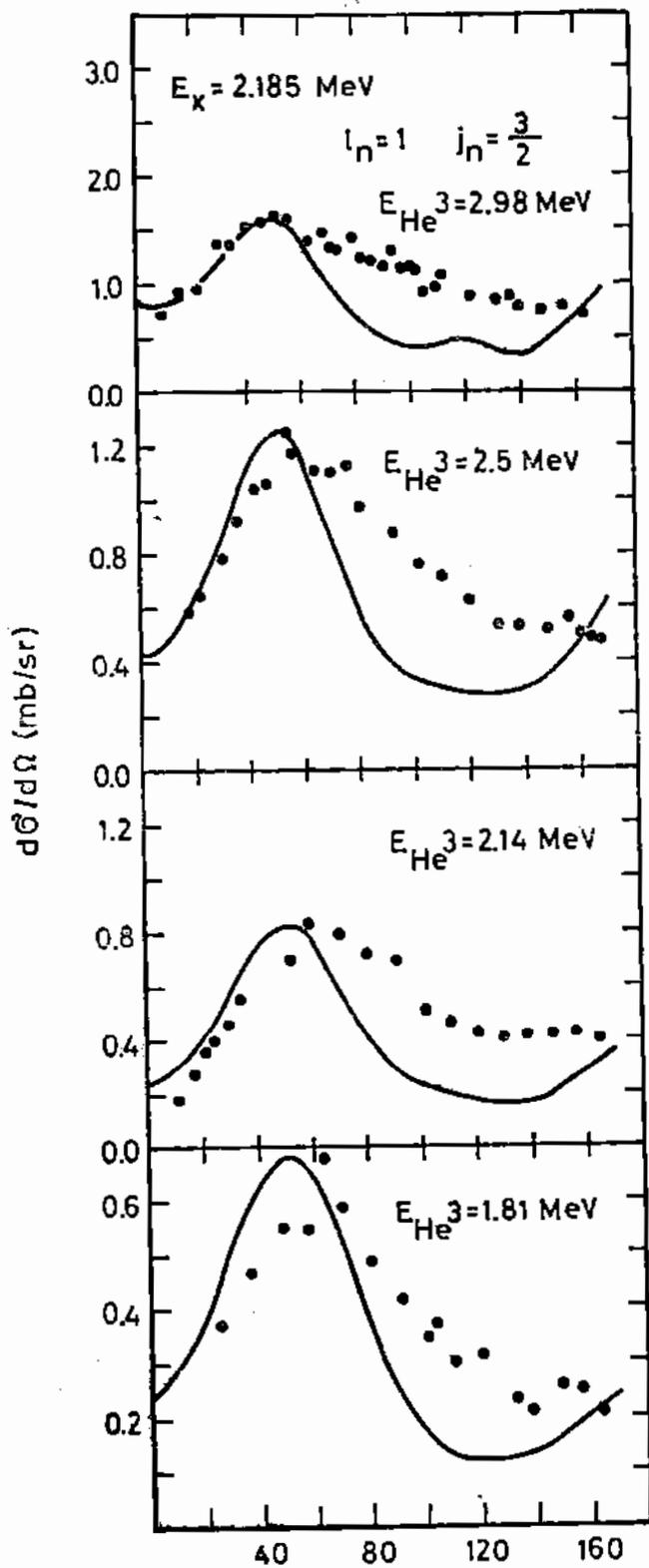
a) α_0 corresponds to the ground state in 6Li with $(J^\pi, \pi) = 1^+, 0$.

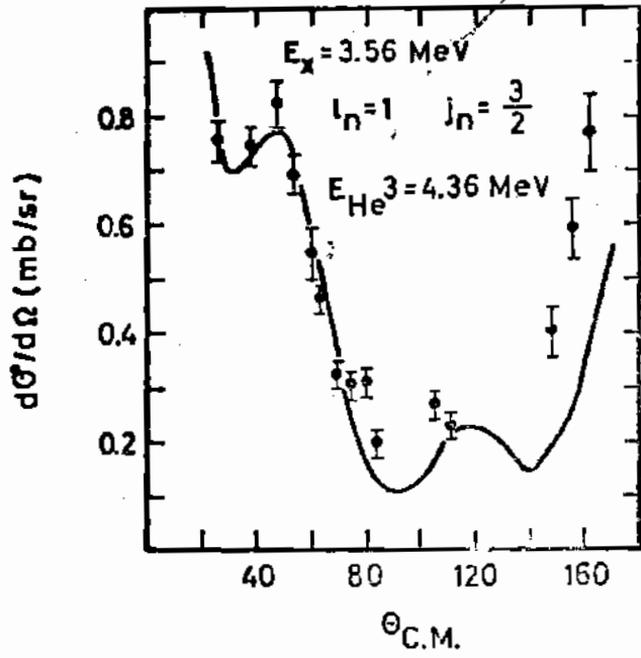
b) α_1 corresponds to the excited state of $E_x = 2.105 MeV$ in 6Li and $(J^\pi, \pi) = 3^+, 0$.

c) α_2 corresponds to the excited state of $E_x = 3.55 MeV$ in 6Li and $(J^\pi, \pi) = 0^+, 1$.

1) $J_n = 1$ for all the three α -groups.







تحليل نتائج التفاعيل ليثيوم^٧ - هيليوم^٣ ه ألنفا^٢ لثيوم^٦
عشر هـ
للثانية من فترة الطاقة ١٨١ الى ١٣٠٥ مليون إلكترون فولت
باعتقاد نظرية الموجه المشروحه .

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الصفحة ٤

في هذا البحث تم تحليل نتائج التوزيع الزاوي لجزئيات الألفا الصادره من
تفاعيل ليثيوم^٧ (هيليوم^٣ ه ألنفا^٢) لثيوم^٦ والمقادير من شحنة
الطاقه ١٨١ الى ١٣٠٥ مليون إلكترون فولت وذلك باستخدام
نظريه الموجه المشروحه . وقد استخدم في ذلك معاملات قوسيه منعه
الليثيوم^٦ والألفا^٢ . وقد تمت النتائج التي تم الحصول عليها من هذا
التحليل بالاحصاءات النظرية .